

# Markov duality for interacting particle systems

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Work based on

- FC, Rouven Frassek and Cristian Giardinà, *Duality for the multispecies stirring process with open boundaries*, 2024 J. Phys. A: Math. Theor. 57 295001.
- Cristian Giardinà and Frank Redig, *Duality for Markov processes: a Lie algebraic approach*, 2025+

# Interacting particle system

- Continuous time Markov process
- Configuration  $\eta(t) = (\eta_1(t), \dots, \eta_L(t))$
- Infinitesimal generator  $\mathcal{L}$

Examples:

- Independent random walker
- Simple Symmetric Exclusion Process

# Markov Duality

Why? "From many to few". Steady state.

## Definition of stochastic duality

Let  $(\eta(t))_{t \geq 0}$  a Markov process with state space  $\Omega$ , and let  $(\xi(t))_{t \geq 0}$  be a Markov process with state space  $\tilde{\Omega}$ . We say that the two process are in a duality relation with duality function  $D : \Omega \times \tilde{\Omega} \rightarrow \mathbb{R}$  if

$$\mathbb{E}^{\eta} [D(\eta(t), \xi)] = \mathbb{E}^{\xi} [D(\eta, \xi(t))] \quad \forall \eta \in \Omega, \forall \xi \in \tilde{\Omega}$$

$\mathcal{L}$  original generator,  $\mathcal{L}^{\text{dual}}$  dual generator

$$(\mathcal{L}D(\cdot, \xi))(\eta) = (\mathcal{L}^{\text{dual}}D(\eta, \cdot))(\xi)$$

Lie algebra generators = elementary building blocks  $\mathcal{L}$  and  $\mathcal{L}^{\text{dual}}$ .

"Find duality between representations of Lie algebras"

# Abstract duality and properties

- $\mathfrak{g}$  Lie algebra
- Pick  $\rho : \mathfrak{g} \rightarrow \mathfrak{gl}(V)$ .  $a \in \mathfrak{g}$  in reps  $\rho$
- Pick  $\hat{\rho} : \mathfrak{g} \rightarrow \mathfrak{gl}(\hat{V})$ .  $\hat{a} \in \mathfrak{g}$  in reps  $\hat{\rho}$ .

$a$  and  $\hat{a}$  duality relation if  $\exists D \in \text{End}(V \otimes \hat{V})$  such that

$$(a \otimes I_{\hat{V}})D = (I_V \otimes \hat{a})D$$

Notation:  $a \xrightarrow{D} \hat{a}$ .

## Properties:

- $a \xrightarrow{D} \hat{a}$ ,  $b \xrightarrow{D} \hat{b}$  then

$$[a, b] \xrightarrow{D} [\hat{b}, \hat{a}]$$

- *New dualities by symmetries:*  $a \xrightarrow{D} \hat{a}$  and  $[a, Q] = 0$ , then  $a \xrightarrow{D'} \hat{a}$  with  $D' = DQ$ .

# Recipe to construct duality

$\mathfrak{g}$  Lie algebra

- $(J_i)_{i \in \{1, \dots, n\}}$  is reps. Let

$$\mathcal{L} = J_1 J_2 \cdots J_m$$

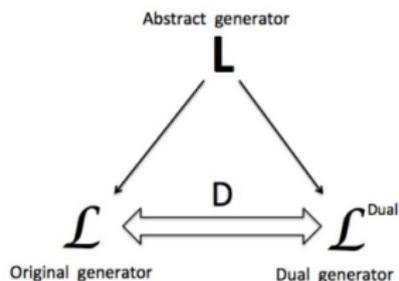
- $(K_i)_{i \in \{1, \dots, n\}}$  conjugate reps;  $d$  such that

$$J_i \xrightarrow{d} K_i \quad \forall i$$

- Dual generator

$$\mathcal{L}^{\text{dual}} = K_m K_{m-1} \cdots K_1$$

- New duality by symmetry



## Example 1: Wright Fisher-Kingman coalescent

Wright-Fisher

$$\mathcal{L} = \frac{1}{2}x(1-x)\frac{d^2}{dx^2}$$

Heisenberg Lie algebra ( $\mathfrak{h}$ , commutators  $[\mathcal{A}, \mathcal{A}^\dagger] = I$ )

$$Af(x) = f'(x), \quad A^\dagger f(x) = xf(x) : [A, A^\dagger] = I$$

$$\mathcal{L} = \frac{1}{2}A^\dagger(I - A^\dagger)A^2$$

Conjugate Lie algebra

$$af(n) = nf(n-1), \quad a^\dagger f(n) = f(n+1) : [a, a^\dagger] = -I$$

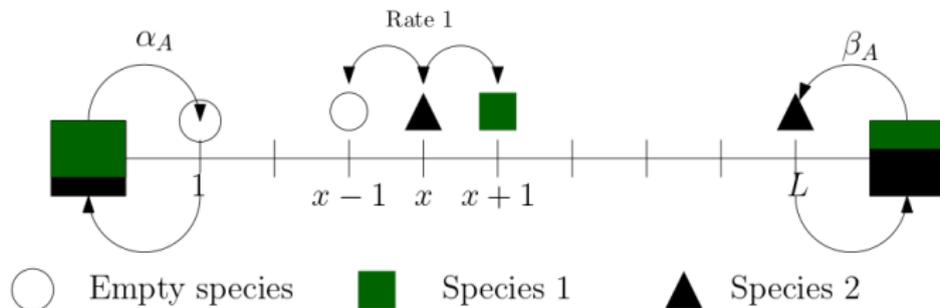
$d(x, n) = x^n$  such that

$$(ad(x, \cdot))(n) = (Ad(\cdot, n))(x) \quad (a^\dagger d(x, \cdot))(n) = (A^\dagger d(\cdot, n))(x)$$

$$\mathcal{L}^{\text{dual}} = \frac{1}{2}a^2(I - a^\dagger)a^\dagger \Rightarrow (\mathcal{L}d(\cdot, n))(x) = (\mathcal{L}^{\text{dual}}d(x, \cdot))(n)$$

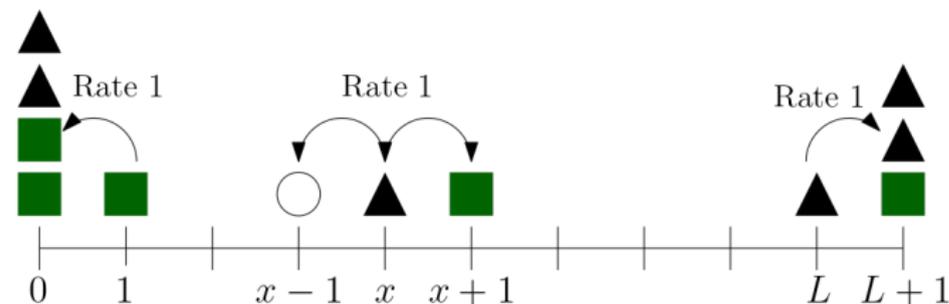
Kingman-Coalescent.

# Multi-species stirring process



$$\mathcal{L} = \mathcal{L}_{\text{left}} + \sum_{x=1}^{L-1} \mathcal{L}_{x,x+1} + \mathcal{L}_{\text{right}}$$

# Absorbing dual process



○ Empty species      ■ Species 1      ▲ Species 2

$$\mathcal{L}^{\text{dual}} = \mathcal{L}_{\text{left}}^{\text{dual}} + \sum_{x=1}^{L-1} \mathcal{L}_{x,x+1} + \mathcal{L}_{\text{right}}^{\text{dual}}$$

## Algebraic description

- Lie algebra  $gl(N)$
- Representation  $\lambda = (1, 0, \dots, 0)$
- Space  $\mathbb{C}^N$  :  $|\eta_1, \dots, \eta_N\rangle$ ,  $\eta_A \in \{0, 1\}$  and  $\eta_N = 1 - \sum_{a=1}^{N-1} \eta_a$
- Algebra generator

$$E_{A,B}|\eta_1, \dots, \eta_N\rangle = \eta_B|\eta_1, \dots, \eta_A + 1, \dots, \eta_B - 1, \dots, \eta_N\rangle$$

Bulk generator (second Casimir!)

$$\mathcal{L}_{x,x+1} = \sum_{A,B=1}^N \left( E_{A,B}^{[x]} E_{B,A}^{[x+1]} - E_{B,B}^{[x]} E_{A,A}^{[x+1]} \right)$$

Boundary generator (left)

$$\mathcal{L}_{\text{left}} = \sum_{A,B=1}^N \alpha_A \left( E_{A,B}^{[1]} - E_{B,B}^{[1]} \right)$$

Duality relation

$$\mathcal{L}^T D = D \mathcal{L}^{\text{dual}} .$$

## Self-duality for the bulk

- take  $E_{A,B} \Rightarrow E_{B,A}$  is conjugate
- $d = \text{Id} \Rightarrow E_{A,B}^T d = d E_{B,A}$
- Dual generator

$$\begin{aligned}\mathcal{L}_{x,x+1}^{\text{dual}} &= \sum_{A,B} E_{B,A}^{[x]} E_{A,B}^{[x+1]} - E_{A,A}^{[x]} E_{B,B}^{[x+1]} \\ &= \mathcal{L}_{x,x+1}\end{aligned}$$

- $\mathcal{L}_{x,x+1}$  central  $\Rightarrow D_x = \exp \left\{ \sum_{a=1}^{N-1} E_{a,N}^{[x]} \right\} d$

satisfy

$$\mathcal{L}_{x,x+1}^T D_x D_{x+1} = D_x D_{x+1} \mathcal{L}_{x,x+1}^{\text{dual}}$$

$$\text{Self duality with } D_{\text{bulk}} = \bigotimes_{x=1}^L D_x$$

## Absorbing dual boundaries

$$D_1^{-1} \mathcal{L}_{\text{left}}^T D_1 = \sum_{a=1}^{N-1} \left( \alpha_a E_{N,a}^{[1]} - E_{a,a}^{[1]} \right)$$

Add extra-site bosonic operator and construct

$$\mathcal{L}_{\text{left}}^{\text{dual}} = \sum_{a=1}^N \left( \left( \mathbf{a}_a^{[0]} \right)^\dagger E_{N,a}^{[1]} - E_{a,a}^{[1]} \right)$$

define extra-site intertwiner  $\mathcal{D}_0$  such that

$$\mathcal{L}_{\text{left}}^T \mathcal{D}_0 D_1 = \mathcal{D}_0 D_1 \mathcal{L}_{\text{left}}^{\text{dual}}$$

Duality matrix

$$D = \mathcal{D}_0 \otimes D_{\text{bulk}} \otimes \mathcal{D}_{L+1}$$

# Exact solution

Integrable model (XXX-quantum chain, extra Yang-Baxter symmetry).

- Duality: triangularize the generator
- Matrix product ansatz (MPA) for the non-equilibrium steady state, i.e. solution to  $H|\Psi\rangle = 0$
- Combining duality with MPA  $\Rightarrow$  exact formulas for  $|\Psi\rangle$

## Currently: non-compact processes

Boundary driven multi-species harmonic process and integrable heat conduction process (IHC)

- Infinite dimensional representations of  $gl(N)$
- Non-trivial symmetries:  $H$  is not the Casimir, but a non-linear function of it.
- Self duality harmonic.
- Duality between harmonic and IHC
- Exact solvability:
  - Physics: Yang-Baxter symmetries, Bethe ansatz solutions
  - Markov processes: Hidden parameter model, propagation of invariant measures
- Kardar-Parisi-Zhang?

# Summary and open questions

## Summary

- Markov duality: study observable via a dual (simpler) process.
- Duality produced via Lie algebra reps
- Contributions:
  - Combine duality and integrable systems to non-equilibrium steady state
  - Extend results to colored processes

## Open question?

- Asymmetric non-compact boundary driven processes
- Mapping non-eq onto eq

Thank you for your attention

## Appendix

$$D(\mathbf{n}, \boldsymbol{\xi}) = \left( \prod_{a=1}^{N-1} \alpha_a^{\xi_a^0} \right) \left( \prod_{x=1}^L \prod_{a=1}^{N-1} \mathbb{1}_{\{n_a^x \geq \xi_a^x\}} \right) \left( \prod_{a=1}^{N-1} \beta_a^{\xi_a^{L+1}} \right)$$

$$|\Psi\rangle = \frac{1}{Z_L} \langle\langle W | \begin{pmatrix} X_1 \\ \vdots \\ X_N \end{pmatrix} \otimes \dots \otimes \begin{pmatrix} X_1 \\ \vdots \\ X_N \end{pmatrix} | V \rangle\rangle$$

$$\lambda = (\nu, 0, \dots, 0)$$

$$D(\mathbf{n}, \boldsymbol{\xi}) = \left( \prod_{a=1}^{N-1} (\alpha_a^x)^{\xi_a^0} \right) \prod_{x=1}^L \left( \frac{(\nu - \sum_{a=1}^{N-1} \xi_a^x)!}{\nu!} \prod_{a=1}^{N-1} \frac{n_a^x!}{(n_a^x - \xi_a^x)!} \right) \left( \prod_{a=1}^{N-1} (\beta_a^x)^{\xi_a^{L+1}} \right)$$